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Structural and spectral properties of a family of deterministic recursive trees: rigorous solutions

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Abstract

As one of the most significant models, the uniform recursive tree (URT) has found many applications in a variety of fields. In this paper, we study rigorously the structural features and spectral properties of the adjacency matrix for a family of deterministic uniform recursive trees (DURTs) that are deterministic versions of URT. Firstly, from the perspective of complex networks, we investigate analytically the main structural characteristics of DURTs, and obtain the accurate solutions for these properties, which include degree distribution, average path length, distribution of node betweenness and degree correlations. Then we determine the complete eigenvalues and their corresponding eigenvectors of the adjacency matrix for DURTs. Our research may shed light on a better understanding of the features for URT. Also, the analytical methods used here are capable of being extended to many other deterministic networks, making the precise computation of their properties (especially the full spectrum characteristics) possible.

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1. Introduction

Structural characterization is very significant for studies in the field of complex networks that have become a focus of attention for the scientific community [1-3]. In the past decade, great efforts have been dedicated to characterizing and understanding the structural properties of real networks [4], including degree distribution, average path length (APL), betweenness, degree correlations, fractality and so forth. These measures have a profound effect on various

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dynamical processes taking place on top of complex networks [5], such as robustness [6–9], epidemic spreading [10, 11], synchronization [12–15] and games [16].

The foregoing measurements focus on direct measurements of structural properties of networks, and play an important role in understanding network complexity [17]. Aside from these measurements there exists a vast literature related to the spectrum of complex networks [18–21], which provides useful insight into the relevant structural properties of and dynamical processes on graphs. In contrast to the fact that structural features capture the static topological properties of complex networks, spectra (eigenvalues and eigenvectors) of adjacency matrix provide global measures of the characterization for network topology. In a variety of dynamical processes, the impact of the network structure is encoded in the spectra of its adjacency matrix, especially the extreme eigenvalues and their corresponding eigenvectors. For example, in the dynamical model for the spreading of infections, the epidemic thresholds are governed by the largest eigenvalue of the adjacency matrix [22, 23], which also plays a fundamental role in determining critical couplings for the onset of coherent behavior [24]. In addition, recent research showed that in the susceptible-infected model of epidemic outbreaks on complex networks, the eigenvectors corresponding to the largest eigenvalue are related to the spreading power of network nodes [25]. In spite of the importance of the eigenvalues and eigenvectors of the adjacency matrix, however, until now, most analysis of the spectra has been confined to approximate or numerical methods, the latter of which is prohibitively time and memory consuming for large-scale networks [18].

On the other hand, in order to mimic real systems and study their structural properties, a great number of network models have been presented [1–3], among which the uniform recursive tree (URT) is perhaps one of the most widely studied models [26]. It is now established that the URT, together with the famous Erdös–Rényi model [27], constitutes the two principal models [28, 29] of random graphs. As one of the simplest trees, the URT is constructed as follows: start with a single node, at each time step, we attach a new node to an existing node selected at random. It has found many important applications in various areas. For example, it has been suggested as models for the spread of epidemics [30], the family trees of preserved copies of ancient or medieval texts [31], chain letter and pyramid schemes [32], to name but a few.

Recently, a class of deterministically growing tree-like networks have been proposed to describe real-world systems whose number of nodes increases exponentially with time [33]. We call them deterministic uniform recursive trees (DURTs), since they are deterministic versions of URT. This kind of deterministic models has received considerable attention from the scientific communities and has turned out to be a useful tool [34–51]. Although uniform recursive tree is well understood [26, 28–31, 52, 53], relatively less is known about the structural and other nature of the DURTs [33].

In this paper, from the viewpoint of complex networks, we offer a comprehensive analysis of the deterministic uniform recursive trees [33]. We firstly determine exactly relevant structural properties of the DURTs, including degree distribution, average path length, betweenness distribution and degree correlations. Then, using methods of graph theory and algebra, we calculate all the eigenvalues and eigenvectors of the adjacency matrix, which are obtained through the recurrence relations derived from the very structure of the DURTs.

2. The deterministic uniform recursive trees

The deterministic uniform recursive trees under consideration are constructed in an iterative way [33]. We denote the trees (networks) after t steps by U_t ($t \ge 0$). Then the networks are built as follows. For t = 0, U_0 is an edge connecting two nodes. For $t \ge 1$, U_t is obtained from



Figure 1. Illustration of a deterministic uniform recursive tree for the special case of m = 1, showing the first several steps of growth process.

 U_{t-1} . We attach *m* new nodes to each node in U_{t-1} . This iterative process is repeated, then we obtain a class of deterministically growing trees with an exponential decreasing spectrum of degrees as shown below. The definition of the model for a particular case of m = 1 is illustrated schematically in figure 1.

We first compute the total number of nodes N_t and the total number of edges E_t in U_t . Let $n_v(t)$ and $n_e(t)$ denote the numbers of nodes and edges created at step t, respectively. Then, $n_v(0) = N_0 = 2$ and $n_e(0) = 1$. By construction, we have $n_v(t) = mN_{t-1}$, thus $N_t = n_v(t) + N_{t-1} = (1+m)N_{t-1}$. Considering the initial condition $N_0 = 2$, we obtain $N_t = 2(1+m)^t$ and $n_v(t) = 2m(1+m)^{t-1}$. Thus, $E_t = N_t - 1 = 2(1+m)^t - 1$. Note that at arbitrary step $t \ge 1$, the addition of each new node leads to only new edge, so $n_e(t) = n_v(t) = 2m(1+m)^{t-1}$ for all $t \ge 1$.

3. Structural properties

In this section, we investigate four important structural properties of U_t , including degree distribution, average path length, betweenness distribution and degree correlations.

3.1. Degree distribution

By definition, the degree of a node *i* is the number of edges connected to *i*. The degree distribution P(k) of a network is the probability that a randomly selected node has exactly *k* edges. Let $k_i(t)$ denote the degree of node *i* at step *t*. If node *i* is added to the network at step t_i , then by construction, $k_i(t_i) = 1$. In each of the subsequent time steps, *m* new nodes will be created connected to *i*. Thus the degree $k_i(t)$ of node *i* satisfies the relation

$$k_i(t) = k_i(t-1) + m.$$
 (1)

Considering the initial condition $k_i(t_i) = 1$, we obtain

$$k_i(t) = 1 + m(t - t_i).$$
 (2)

Since the degree of each node has been obtained explicitly as in equation (2), we can get the degree distribution via its cumulative distribution [3]

$$P_{\rm cum}(k) = \sum_{k'=k}^{\infty} P(k'), \tag{3}$$

which is the probability that the degree is greater than or equal to k. An important advantage of the cumulative distribution is that it can reduce the noise in the tail of probability distribution. Moreover, for some networks whose degree distributions have exponential tails: $P(\tilde{k}) \sim e^{-\tilde{k}/\kappa}$, the cumulative distribution also has an exponential expression with the same exponent:

$$P_{\rm cum}(\tilde{k}) = \sum_{k'=\tilde{k}}^{\infty} P(k') \sim \sum_{k'=\tilde{k}}^{\infty} e^{-k'/\kappa} \sim e^{-\tilde{k}/\kappa}.$$
(4)

This makes exponential distributions particularly easy to detect experimentally, by plotting the corresponding cumulative distributions on semilogarithmic scales.

Using equation (2), we have $P_{\text{cum}}(k) = \sum_{k'=k}^{\infty} P(k') = P(t' \leq \tau = t - \frac{k-1}{m})$. Hence

$$P_{\rm cum}(k) = \sum_{t'=0}^{\tau} \frac{n_v(t')}{N_t} = \frac{2(1+m)^{t-\frac{k-1}{m}}}{2(1+m)^t} = (1+m)^{-\frac{k-1}{m}},\tag{5}$$

which decays exponentially with k. Note that when m > 1, the possible degrees are not arbitrary, equation (5) holds only for those k being equal to 1 modulo m. Thus the DURTs are a family of exponential networks, which have a similar form of degree distribution as its stochastic version—the URT [29].

3.2. Average path length

Average path length means the minimum number of edges connecting a pair of nodes, averaged over all node pairs. It is defined to be

$$\bar{d}_t = \frac{S_t}{N_t (N_t - 1)/2},$$
(6)

where S_t denotes the sum of the total distances between two nodes over all pairs, that is

$$S_t = \sum_{i \neq j} d_{i,j},\tag{7}$$

where $d_{i,j}$ is the shortest distance between node *i* and *j*. Note that in equation (7), for a couple of nodes *i* and *j* ($i \neq j$), we only count $d_{i,j}$ or $d_{j,i}$, not both.

Let Ω_{new}^t and Ω_{old}^t represent the sets of nodes created at step *t* or earlier, respectively. Then one can write the sum over all shortest paths S_t in network U_t as

$$S_t = \sum_{i \in \Omega_{\text{new}}^t, j \in \Omega_{\text{old}}^t} d_{i,j} + \sum_{i \in \Omega_{\text{new}}^t, j \in \Omega_{\text{new}}^t} d_{i,j} + \sum_{i \in \Omega_{\text{old}}^t, j \in \Omega_{\text{old}}^t} d_{i,j},$$
(8)

where the third term is exactly S_{t-1} , i.e.,

$$\sum_{i \in \Omega_{\text{old}}^t, j \in \Omega_{\text{old}}^t} d_{i,j} = S_{t-1}.$$
(9)

By construction, we can obtain the following relations for the first and second terms on the right-hand side of equation (8):

$$\sum_{i \in \Omega_{\text{new}}^{t}, j \in \Omega_{\text{old}}^{t}} d_{i,j} = m \left(N_{t-1}^{2} + 2S_{t-1} \right), \tag{10}$$

$$\sum_{i \in \Omega_{new}^{t}, j \in \Omega_{new}^{t}} d_{i,j} = m^{2} S_{t-1} + m N_{t-1} (m N_{t-1} - 1).$$
(11)

The term $mN_{t-1}(mN_{t-1}-1)$ in equation (11) pops out from counting: each path connecting two new points comes from a path connecting two old points by adding two edges, is by

increasing the length by 2. As there are $\frac{1}{2}mN_{t-1}(mN_{t-1}-1)$ pairs of new points, the total increase in length is $mN_{t-1}(mN_{t-1}-1)$.

Substituting equations (9)–(11) into equation (8) and considering $N_t = 2(1 + m)^t$, the total distance is obtained to be

$$S_{t} = (1+m)^{2} S_{t-1} + m(1+m) N_{t-1}^{2} - m N_{t-1}$$

= $(1+m)^{2t} S_{0} + m(1+m) \sum_{i=0}^{t-1} (1+m)^{2(t-1-i)} N_{i}^{2} - m \sum_{i=0}^{t-1} (1+m)^{2(t-1-i)} N_{i}$
= $(4mt+m-1)(1+m)^{2t-1} + 2(1+m)^{t-1}$. (12)

Inserting equation (12) into (6), we have

$$\bar{d}_{t} = \frac{2[(4mt+m-1)(1+m)^{2t-1}+2(1+m)^{t-1}]}{2(1+m)^{t}[2(1+m)^{t}-1]}$$
$$= \frac{(1+m)^{t}(4mt+m-1)+2}{2(1+m)^{t+1}-(1+m)}.$$
(13)

In the infinite network size limit $(t \to \infty)$,

$$\bar{d}_t \cong \frac{2m}{m+1}t + \frac{m-1}{2(m+1)} = \frac{\ln N_t}{\ln(m+1)} - \frac{\ln 2}{\ln(m+1)} + \frac{m-1}{2(m+1)},$$
(14)

which means that the average path length shows a logarithmic scaling with the size of the network, indicating a similar small-world behavior as the URT [53] and the Watts–Strogatz (WS) model [54].

3.3. Betweenness distribution

Betweenness of a node is the accumulated fraction of the total number of shortest paths going through the given node over all node pairs [55, 56]. More precisely, the betweenness of a node i is

$$b_i = \sum_{j \neq i \neq k} \frac{\sigma_{jk}(i)}{\sigma_{jk}},\tag{15}$$

where σ_{jk} is the total number of the shortest path between node j and k, and $\sigma_{jk}(i)$ is the number of the shortest path running through node i.

Since for a tree, there is a unique shortest path between each pair of nodes [57–61]. Thus the betweenness of a node is simply given by the number of distinct shortest paths passing through the node. Then in U_t , the betweenness of a τ -generation-old node v, which is created at step $t - \tau + 1$, denoted as $b_t(\tau)$ becomes

$$b_t(\tau) = \Theta_t^{\tau} \left[N_t - \left(\Theta_t^{\tau} + 1 \right) \right] + \frac{\Theta_t^{\tau} (\Theta_t^{\tau} - 1)}{2} - \sum_{h=2}^{\tau-1} m \frac{\Theta_t^h \left(\Theta_t^h + 1 \right)}{2}, \tag{16}$$

where Θ_t^{τ} denotes the total number of descendants of node *v* at time *t*, where the descendants of a node are its children, its children's children and so on. Note that the descendants of node *v* exclude *v* itself. The first term in equation (16) counts the shortest paths from descendants of *v* to other vertices. The second term accounts for the shortest paths between descendants of *v*. The third term describes the shortest paths between descendants of *v* that do not pass through *v*.

To find $b_t(\tau)$, it is necessary to explicitly determine the descendants Θ_t^{τ} of node v, which is related to that of v's children via [59, 60]

$$\Theta_t^{\tau} = \sum_{j=1}^{\tau-1} m \big(\Theta_t^j + 1 \big). \tag{17}$$

Using $\Theta_t^1 = 0$, we can solve equation (17) inductively,

$$\Theta_t^{\tau} = (1+m)^{\tau-1} - 1.$$
(18)

Substituting the result of equation (18) and $N_t = 2(1 + m)^t$ into equation (16), we have

$$b_t(\tau) = 2(1+m)^{t+\tau-1} - 2(1+m)^t - \frac{(m+3)[(1+m)^{2(\tau-1)} - 1]}{2(m+2)},$$
(19)

which is approximately equal to $2(m + 1)^{t+\tau-1}$ for large τ . Then the cumulative betweenness distribution is

$$P_{\text{cum}}(b) = \sum_{\mu \leqslant t - \tau + 1} \frac{n_v(\mu)}{N_t}$$

= $\frac{(1+m)^{t+1}}{(1+m)^{t+\tau}} \approx \frac{N_t}{b} \sim b^{-1},$ (20)

which shows that the betweenness distribution exhibits a power-law behavior with exponent $\gamma_b = 2$, the same scaling has been also obtained for the URT [52] and the m = 1 case of the Barabási–Albert (BA) model [62] describing a random scale-free treelike network [57, 58]. Therefore, power-law betweenness distribution is not an exclusive property of scale-free networks.

3.4. Degree correlations

An interesting quantity related to degree correlations [63] is the average degree of the nearest neighbors for nodes with degree k, denoted as $k_{nn}(k)$ [64–66]. When $k_{nn}(k)$ increases with k, it means that nodes have a tendency to connect to nodes with a similar or larger degree. In this case the network is defined as assortative [67]. In contrast, if $k_{nn}(k)$ is decreasing with k, which implies that nodes of large degree are likely to have near neighbors with small degree, then the network is said to be disassortative. If correlations are absent, $k_{nn}(k) = \text{const.}$

For U_t , we can exactly calculate $k_{nn}(k)$. Except for the initial two nodes generated at step 0, no nodes born at the same step, which have the same degree, will be linked to each other. All links to nodes with larger degree are made at the creation step, and then links to nodes with smaller degree are made at each subsequent step. This results in the expression for $k = 1 + m(t - t_i)$ ($t_i \ge 1$)

$$k_{\rm nn}(k) = \frac{1}{n_v(t_i)k(t_i, t)} \left[\sum_{t_i'=0}^{t_i'=t_i-1} mn_v(t_i')k(t_i', t) + \sum_{t_i'=t_i+1}^{t_i'=t} mn_v(t_i)k(t_i', t) \right], \quad (21)$$

where $k(t_i, t)$ represents the degree of a node at step t, which was generated at step t_i . Here the first sum on the right-hand side accounts for the links made to nodes with larger degree (i.e. $t'_i < t_i$) when the node was generated at t_i . The second sum describes the links made to the current smallest degree nodes at each step $t'_i > t_i$.

After some algebraic manipulations, equation (21) is simplified to

$$k_{nn}(k) = \frac{2mt + 2 - 2mt_i + m}{1 + m(t - t_i)} + \frac{m^2(t - t_i - 1)(t - t_i)}{2[1 + m(t - t_i)]} - \frac{(1 + m)^{1 - t_i}}{1 + m(t - t_i)}.$$
(22)

Writing equation (22) in terms of k, it is straightforward to obtain

$$k_{\rm nn}(k) = \frac{k}{2} + \frac{2-m}{2} + \frac{3m+1}{2k} - \frac{(1+m)^{1+\frac{k-1}{m}}}{k \cdot (1+m)^t}.$$
(23)

Thus we have obtained the degree correlations for those nodes born at $t_i \ge 1$. For the initial two nodes, each has a degree of k = 1 + mt, and it is easy to obtain

$$k_{nn}(k = 1 + mt) = \frac{1}{k} \left(\sum_{t'_i=1}^{t'_i=t} mk(t'_i, t) + k(0, t) \right)$$
$$= \frac{k}{2} + \frac{2 - m}{2} + \frac{m - 1}{2k}.$$
(24)

From equations (23) and (24), it is obvious that for a large network (i.e. $t \to \infty$), $k_{nn}(k)$ is approximately a linear function of k, which shows that the network is assortative.

4. Eigenvalues and eigenvectors of the adjacency matrix

As known from section 2, there are $2(m + 1)^t$ vertices in U_t . We denote by V_t the vertex set of U_t , i.e., $V_t = \{v_1, v_2, \dots, v_{2(m+1)'}\}$. Let $\mathbf{A}_t = [a_{ij}]$ be the adjacency matrix of network U_t , where $a_{ij} = a_{ji} = 1$ if nodes *i* and *j* are connected, $a_{ij} = a_{ji} = 0$ otherwise. For an arbitrary graph, it is generally difficult to determine all eigenvalues and the corresponding eigenvectors of its adjacency matrix, but below we will show that for U_t one can settle this problem.

4.1. Eigenvalues

We begin by studying the eigenvalues of U_t . By construction, it is easy to find that the adjacency matrix A_t satisfies the following relation:

$$\mathbf{A}_{t} = \begin{pmatrix} \mathbf{A}_{t-1} & \mathbf{I} & \mathbf{I} & \cdots & \mathbf{I} \\ \mathbf{I} & \mathbf{0} & \mathbf{0} & \cdots & \mathbf{0} \\ \mathbf{I} & \mathbf{0} & \mathbf{0} & \cdots & \mathbf{0} \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ \mathbf{I} & \mathbf{0} & \mathbf{0} & \cdots & \mathbf{0} \end{pmatrix}_{(m+1) \times (m+1)}$$
(25)

where each block is a $2(m+1)^{t-1} \times 2(m+1)^{t-1}$ matrix and **I** is the identity matrix. Then, the characteristic polynomial of A_t is

$$P_{t}(x) = \det(x\mathbf{I} - \mathbf{A}_{t})$$

$$= \det\begin{pmatrix} x\mathbf{I} - \mathbf{A}_{t-1} & -\mathbf{I} & -\mathbf{I} & \cdots & -\mathbf{I} \\ -\mathbf{I} & x\mathbf{I} & \mathbf{0} & \cdots & \mathbf{0} \\ -\mathbf{I} & \mathbf{0} & x\mathbf{I} & \cdots & \mathbf{0} \\ \vdots & \vdots & \vdots & \vdots \\ -\mathbf{I} & \mathbf{0} & \mathbf{0} & \cdots & x\mathbf{I} \end{pmatrix}$$

$$= (\det(x\mathbf{I}))^{m} \cdot \det\begin{pmatrix} (x - \frac{m}{x})\mathbf{I} - \mathbf{A}_{t-1} & \mathbf{0} & \mathbf{0} & \cdots & \mathbf{0} \\ -\frac{1}{x}\mathbf{I} & \mathbf{I} & \mathbf{0} & \cdots & \mathbf{0} \\ -\frac{1}{x}\mathbf{I} & \mathbf{0} & \mathbf{I} & \cdots & \mathbf{0} \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ -\frac{1}{x}\mathbf{I} & \mathbf{0} & \mathbf{0} & \cdots & \mathbf{I} \end{pmatrix}, \quad (26)$$

where the elementary operations of matrix have been used. According to the results in [68], we have

$$P_t(x) = (\det(x\mathbf{I}))^m \cdot \det\left(\left(x - \frac{m}{x}\right)\mathbf{I} - \mathbf{A}_{t-1}\right).$$
(27)

Thus, $P_t(x)$ can be written recursively as follows:

$$P_t(x) = x^{2m(m+1)^{t-1}} \cdot P_{t-1}(\varphi(x)),$$
(28)

where $\varphi(x) = x - \frac{m}{x}$. This recursive relation given by equation (28) is very important, from which we will determine the complete eigenvalues of U_t and their corresponding eigenvectors. Note that $P_{t-1}(x)$ is a monic polynomial of degree $2(m + 1)^{t-1}$, then the exponent of $\frac{m}{x}$ in $P_{t-1}(\varphi(x))$ is $2(m + 1)^{t-1}$, and hence the exponent of factor x in $P_t(x)$ is

$$2m(m+1)^{t-1} - 2(m+1)^{t-1} = 2(m-1)(m+1)^{t-1}.$$
(29)

Consequently, 0 is an eigenvalue of A_t , and its multiplicity is $2(m-1)(m+1)^{t-1}$.

Note that U_t has $2(m + 1)^t$ eigenvalues. We represent these $2(m + 1)^t$ eigenvalues as $\lambda_1^t, \lambda_2^t, \ldots, \lambda_{2(m+1)^t}^t$, respectively. For convenience, we presume $\lambda_1^t \leq \lambda_2^t \leq \cdots \leq \lambda_{2(m+1)^t}^t$, and denote by AE_t the set of these eigenvalues, i.e. $AE_t = \{\lambda_1^t, \lambda_2^t, \ldots, \lambda_{2(m+1)^t}^t\}$. All the eigenvalues in set AE_t can be divided into two parts. According to the above analysis, $\lambda = 0$ is an eigenvalue with multiplicity $2(m-1)(m+1)^{t-1}$, which provide parts of the eigenvalues of A_t . We denote by AE_t^t the set of eigenvalues 0 of U_t , i.e.

$$AE'_{t} = \{\underbrace{0, 0, 0, \dots, 0, 0}_{2(m-1)(m+1)^{t-1}}\}$$
(30)

It should be noted that here we neglect the distinctness of elements in the set. The remaining $4(m+1)^{t-1}$ adjacency eigenvalues of U_t are determined by the equation $P_{t-1}(\varphi(x)) = 0$. Let these $4(m+1)^{t-1}$ eigenvalues be $\tilde{\lambda}_1^t, \tilde{\lambda}_2^t, \ldots, \tilde{\lambda}_{4(m+1)^{t-1}}^t$, respectively. For convenience, we presume $\tilde{\lambda}_1^t \leq \tilde{\lambda}_2^t \leq \cdots \leq \tilde{\lambda}_{4(m+1)^{t-1}}^t$, and denote by AE_t^* the set of these eigenvalues, i.e. $AE_t^* = {\tilde{\lambda}_1^t, \tilde{\lambda}_2^t, \ldots, \tilde{\lambda}_{4(m+1)^{t-1}}^t}$. Therefore, the eigenvalue set of U_t can be expressed as $AE_t = AE_t^t \cup AE_t^*$.

From equation (28), we have that for an arbitrary element in AE_{t-1} , say $\lambda_i^{t-1} \in AE_{t-1}$, both solutions of $x - \frac{m}{x} = \lambda_i^{t-1}$ are in AE_t^* . In fact, equation $x - \frac{m}{x} = \lambda_i^{t-1}$ is equivalent to

$$x^2 - \lambda_i^{t-1} x - m = 0. (31)$$

We use notations $\tilde{\lambda}_i^t$ and $\tilde{\lambda}_{i+2(m+1)^{t-1}}^t$ to represent the two solutions to equation (31), since they provide a natural increasing order of the eigenvalues of U_t , which can be seen from below argument. Solving this quadratic equation, its roots are obtained to be $\tilde{\lambda}_i^t = r_1(\lambda_i^{t-1})$ and $\tilde{\lambda}_{i+2(m+1)^{t-1}}^t = r_2(\lambda_i^{t-1})$, where the functions $r_1(\lambda)$ and $r_2(\lambda)$ satisfy

$$r_1(\lambda) = \frac{1}{2} \left(\lambda - \sqrt{\lambda^2 + 4m} \right), \tag{32}$$

$$r_2(\lambda) = \frac{1}{2} \left(\lambda + \sqrt{\lambda^2 + 4m} \right). \tag{33}$$

Substituting each adjacency eigenvalue of U_{t-1} into equations (32) and (33), we can obtain the set AE_t^* of eigenvalues of U_t . Since $AE_0 = \{-1, 1\}$, by recursively applying the functions provided by equations (32) and (33), the eigenvalues of U_t can be determined completely.

It is obvious that both $r_1(\lambda)$ and $r_2(\lambda)$ are monotonously increasing functions. On the other hand, since $r_1(\lambda) = \frac{1}{2} \left(\lambda - \sqrt{\lambda^2 + 4m} \right) = \frac{-2m}{(\lambda + \sqrt{\lambda^2 + 4m})}$, so $r_1(\lambda) < 0$. Similarly, we can show that $r_2(\lambda) > 0$. Thus for arbitrary fixed $\lambda', r_1(\lambda) < r_2(\lambda')$ holds for all λ . Then we have the



Figure 2. The distribution of eigenvalues $P_t(\lambda)$ defined as the ratio between the multiplicity of eigenvalue λ and the network order N_t .

following conclusion: if the eigenvalues set of U_{t-1} is $AE_{t-1} = \{\lambda_1^{t-1}, \lambda_2^{t-1}, \dots, \lambda_{2(m+1)^{t-1}}^{t-1}\}$, then solving equations (32) and (33) we can obtain the eigenvalue set AE_t^* of U_t to be $AE_t^* = \{\tilde{\lambda}_1^t, \tilde{\lambda}_2^t, \dots, \tilde{\lambda}_{4(m+1)^{t-1}}^t\}$, where $\tilde{\lambda}_1^t \leq \tilde{\lambda}_2^t \leq \dots \leq \tilde{\lambda}_{2(m+1)^{t-1}}^t < 0 < \tilde{\lambda}_{2(m+1)^{t-1}+1}^t \leq \tilde{\lambda}_{2(m+1)^{t-1}+2}^t \leq \dots \leq \tilde{\lambda}_{4(m+1)^{t-1}}^t$. Recall that AE_t' consists of $2(m-1)(m+1)^{t-1}$ elements, all of which are 0, so we can easily get the eigenvalue set of U_t to be $AE_t = AE_t^* \cup AE_t'$.

For above arguments, we can easily see that for the case of m = 1, all the 2^{t+1} eigenvalues of U_t are different, which is an interesting feature and has less been previously reported in other network models thus may have some far-reaching consequences. For other m > 1, some eigenvalues multiple. In figure 2 we plot the distribution of eigenvalues for two cases: m = 2and m = 3. It is observed that different from the uniform distribution of m = 1 case, for m > 1, the distribution of eigenvalues exhibits the form of peaks.

4.2. Eigenvectors

Similar to the eigenvalues, the eigenvectors of \mathbf{A}_t follow directly from those of \mathbf{A}_{t-1} . Assume that λ is an arbitrary eigenvalue of U_t , whose corresponding eigenvector is $\boldsymbol{v} \in \mathbf{R}^{2(m+1)^t}$, where $\mathbf{R}^{2(m+1)^t}$ represents the $2(m+1)^t$ -dimensional vector space. Then we can solve equation $(\lambda \mathbf{I} - \mathbf{A}_t)\boldsymbol{v} = 0$ to find the eigenvector \boldsymbol{v} . We distinguish two cases: $\lambda \in AE_t^*$ and $\lambda \in AE_t'$, which will be addressed in detail as follows.

For the first case $\lambda \in AE_t^*$, we can rewrite the equation $(\lambda \mathbf{I} - \mathbf{A}_t)\mathbf{v} = 0$ as

$$\begin{pmatrix} \lambda \mathbf{I} - \mathbf{A}_{t-1} & -\mathbf{I} & -\mathbf{I} & \cdots & -\mathbf{I} \\ -\mathbf{I} & \lambda \mathbf{I} & \mathbf{0} & \cdots & \mathbf{0} \\ -\mathbf{I} & \mathbf{0} & \lambda \mathbf{I} & \cdots & \mathbf{0} \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ -\mathbf{I} & \mathbf{0} & \mathbf{0} & \cdots & \lambda \mathbf{I} \end{pmatrix} \begin{pmatrix} \boldsymbol{v}_1 \\ \boldsymbol{v}_2 \\ \boldsymbol{v}_3 \\ \vdots \\ \boldsymbol{v}_{m+1} \end{pmatrix} = 0,$$
(34)

where vector v_i $(1 \le i \le m + 1)$ are components of v. Equation (34) leads to the following equations:

$$(\lambda \mathbf{I}_{t-1} - \mathbf{A}_{t-1}) v_1 - v_2 - \dots - v_{m+1} = \mathbf{0},$$
(35)

$$-\boldsymbol{v}_1 + \lambda \boldsymbol{v}_i = \boldsymbol{0} \qquad (2 \leqslant i \leqslant m+1). \tag{36}$$

Resolve equation (36) to find

$$\boldsymbol{v}_i = \frac{1}{\lambda} \boldsymbol{v}_1 \qquad (2 \leqslant i \leqslant m+1). \tag{37}$$

Substituting equation (37) into (35) we have

$$\left[\left(\lambda - \frac{m}{\lambda}\right)\mathbf{I} - \mathbf{A}_{t-1}\right]\boldsymbol{v}_1 = 0, \tag{38}$$

which indicates that v_1 is the solution to equation (35) while v_i ($2 \le i \le m+1$) are uniquely decided by v_1 via equation (37).

From equation (28) in the preceding subsection, it is clear that if λ is an eigenvalue of adjacency matrix \mathbf{A}_t , then $f(\lambda) = \lambda - \frac{m}{\lambda}$ must be one eigenvalue of \mathbf{A}_{t-1} . (Recall that if $\lambda = \tilde{\lambda}_i^t \in AE_t^*$, then $\varphi(\tilde{\lambda}_i^t) = \lambda_i^{t-1}$ for $i \leq 2(m+1)^{t-1}$, or $\varphi(\tilde{\lambda}_i^t) = \lambda_{i-2(m+1)^{t-1}}$ for $i > 2(m+1)^{t-1}$). Thus, equation (38) together with equation (28) shows that v_1 is an eigenvector of matrix \mathbf{A}_{t-1} corresponding to the eigenvalue $\lambda - \frac{m}{\lambda}$ determined by λ , while

$$\boldsymbol{v} = \begin{pmatrix} \boldsymbol{v}_1 \\ \boldsymbol{v}_2 \\ \boldsymbol{v}_3 \\ \vdots \\ \boldsymbol{v}_{m+1} \end{pmatrix} = \begin{pmatrix} \boldsymbol{v}_1 \\ \frac{1}{\lambda} \boldsymbol{v}_1 \\ \frac{1}{\lambda} \boldsymbol{v}_1 \\ \vdots \\ \frac{1}{\lambda} \boldsymbol{v}_1 \end{pmatrix}$$
(39)

is an eigenvector of \mathbf{A}_t corresponding to the eigenvalue λ .

Since for the initial graph U_0 , its adjacency matrix \mathbf{A}_0 has two eigenvalues -1 and 1 with respective eigenvectors $(1, -1)^{\top}$ and $(1, 1)^{\top}$. By recursively applying the above process, we can obtain all the eigenvectors corresponding to $\lambda \in AE_t^*$.

For the second case of $\lambda \in AE'_t$, where all $\lambda = 0$, the equation $(\lambda \mathbf{I} - \mathbf{A}_t)\mathbf{v} = 0$ can be recast as

$$\begin{pmatrix} -\mathbf{A}_{t-1} & -\mathbf{I} & -\mathbf{I} & \cdots & -\mathbf{I} \\ -\mathbf{I} & \mathbf{0} & \mathbf{0} & \cdots & \mathbf{0} \\ -\mathbf{I} & \mathbf{0} & \mathbf{0} & \cdots & \mathbf{0} \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ -\mathbf{I} & \mathbf{0} & \mathbf{0} & \cdots & \mathbf{0} \end{pmatrix} \begin{pmatrix} \mathbf{v}_1 \\ \mathbf{v}_2 \\ \mathbf{v}_3 \\ \vdots \\ \mathbf{v}_{m+1} \end{pmatrix} = \mathbf{0},$$
(40)

where vector v_i $(1 \le i \le m + 1)$ are components of v. Equation (40) leads to the following equations:

$$\boldsymbol{v}_1 = \boldsymbol{0},\tag{41}$$

$$v_2 + v_3 + \dots + v_{m+1} = \mathbf{0}.$$
 (42)

From equation (41), v_1 is a zero vector, and we denote by $v_{i,j}$ the *j* th component of the column vector v_i . Equation (42) gives us the following equations:

$$\begin{cases} \boldsymbol{v}_{2,1} + \boldsymbol{v}_{3,1} + \cdots + \boldsymbol{v}_{m+1,1} = \boldsymbol{0} \\ \boldsymbol{v}_{2,2} + \boldsymbol{v}_{3,2} + \cdots + \boldsymbol{v}_{m+1,2} = \boldsymbol{0} \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ \boldsymbol{v}_{2,2(m+1)^{t-1}} + \boldsymbol{v}_{3,2(m+1)^{t-1}} + \cdots + \boldsymbol{v}_{m+1,2(m+1)^{t-1}} = \boldsymbol{0}. \end{cases}$$

The set of all solutions to any equation above consists of vectors that can be written as

$$\begin{pmatrix} \boldsymbol{v}_{2,j} \\ \boldsymbol{v}_{3,j} \\ \boldsymbol{v}_{4,j} \\ \vdots \\ \boldsymbol{v}_{m+1,j} \end{pmatrix} = k_{1,j} \begin{pmatrix} -1 \\ 1 \\ 0 \\ \vdots \\ 0 \end{pmatrix} + k_{2,j} \begin{pmatrix} -1 \\ 0 \\ 1 \\ \vdots \\ 0 \end{pmatrix} + \dots + k_{m-1,j} \begin{pmatrix} -1 \\ 0 \\ 0 \\ \vdots \\ 1 \end{pmatrix},$$
(43)

where $k_{1,j}, k_{2,j}, \ldots, k_{m-1,j}$ are any real numbers. From equation (43), the solutions for all the vectors v_i ($2 \le i \le m+1$) can be rewritten as

$$\begin{pmatrix} \boldsymbol{v}_{2}^{\vee} \\ \boldsymbol{v}_{3}^{\top} \\ \boldsymbol{v}_{4}^{\vee} \\ \vdots \\ \boldsymbol{v}_{m+1}^{\vee} \end{pmatrix} = \begin{pmatrix} -1 & -1 & \cdots & -1 \\ 1 & \mathbf{0} & \cdots & \mathbf{0} \\ \mathbf{0} & \mathbf{1} & \cdots & \mathbf{0} \\ \vdots & \vdots & & \vdots \\ \mathbf{0} & \mathbf{0} & \cdots & \mathbf{1} \end{pmatrix} \qquad \begin{pmatrix} k_{1,1} & k_{1,2} & \cdots & k_{1,2(m+1)^{t-1}} \\ k_{2,1} & k_{2,2} & \cdots & k_{2,2(m+1)^{t-1}} \\ k_{3,1} & k_{3,2} & \cdots & k_{3,2(m+1)^{t-1}} \\ \vdots & \vdots & & \vdots \\ k_{m-1,1} & k_{m-1,2} & \cdots & k_{m-1,2(m+1)^{t-1}} \end{pmatrix}, \quad (44)$$

where $k_{i,j}$ $(1 \le i \le m - 1; 1 \le j \le 2(m + 1)^{t-1}))$ are arbitrary real numbers. According to equation (44), we can obtain the eigenvector v corresponding to the eigenvalue **0**. Moreover, it is easy to see that the dimension of the eigenspace of matrix A_t associated with eigenvalue **0** is $2(m - 1)(m + 1)^{t-1}$.

5. Conclusion and discussion

In conclusion, we have studied a family of deterministic models for the uniform recursive tree, which we name the deterministic uniform recursive trees that are constructed in a recursive way. The DURTs are in fact deterministic variants of the intensively studied random uniform recursive tree. We have presented an exhaustive analysis of various structural properties of the DURTs, and obtained the precise solutions for these features that include degree distributions, average path length, betweenness distribution and degree correlations. Aside from their deterministic structures, the obtained structural characteristics of the DURTs are similar to those of the URT. Consequently, the DURTs may provide useful insight to the practices as the URT.

Furthermore, by using the methods of linear algebra and graph theory, we have performed a detailed analysis of the complete eigenvalues and their corresponding eigenvectors of the adjacency matrix for the DURTs. We have fully characterized the spectral properties and eigenvectors for the DURTs. We have shown that all the eigenvalues and eigenvectors of the adjacency matrix for the DURTs can be directly determined from those for the initial network. It is expected that the methods applied here can be extended to a larger type of deterministic networks.

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References

- [1] Albert R and Barabási A-L 2002 Rev. Mod. Phys. 74 47
- [2] Dorogvtsev S N and Mendes J F F 2002 Adv. Phys. 51 1079
- [3] Newman M E J 2003 SIAM Rev. 45 167
- [4] da F Costa L, Rodrigues F A, Travieso G and Boas P R V 2007 Adv. Phys. 56 167
- [5] Dorogovtsev S N, Goltsev A V and Mendes J F F 2008 Rev. Mod. Phys. 80 1275
- [6] Albert R, Jeong H and Barabási A-L 2000 Nature 406 378
- [7] Callaway D S, Newman M E J, Strogatz S H and Watts D J 2000 Phys. Rev. Lett. 85 5468
- [8] Cohen R, Erez K, ben-Avraham D and Havlin S 2001 Phys. Rev. Lett. 86 3682
- [9] Zhang Z Z, Zhou S G and Zou T 2007 Eur. Phys. J. B 56 259
- [10] Pastor-Satorras R and Vespignani A 2001 Phys. Rev. Lett. 86 3200
- [11] Zhang Z Z, Zhou S G, Zou T and Chen G S 2008 J. Stat. Mech.: Theory Exp. P09008
- [12] Barahona M and Pecora L M 2002 Phys. Rev. Lett. 89 054101
- [13] Zhang Z Z, Rong L L and Zhou S G 2006 Phys. Rev. E 74 046105
- [14] Comellas F and Gago S 2007 J. Phys. A: Math. Theor. 40 4483
- [15] Arenas A, Díaz-Guilera A, Kurths J, Moreno Y and Zhou C S 2008 Phy. Rep. 469 93
- [16] Szabó G and Fáth G 2007 Phy. Rep. 446 97
- [17] Barabási A-L 2005 Nat. Phys. 1 68
- [18] Farkas I J, Derényi I, Barabási A-L and Vicsek T 2001 Phys. Rev. E 64 026704
- [19] Goh K-I, Kahng B and Kim D 2001 Phys. Rev. E 64 051903
- [20] Dorogovtsev S N, Goltsev A V, Mendes J F F and Samukhin A N 2003 Phys. Rev. E 68 046109
- [21] Chung F, Lu L and Vu V 2003 Proc. Natl. Acad. Sci. USA 100 6313
- [22] Boguñá M, Pastor-Satorras R and Vespignani A 2003 Phys. Rev. Lett. 90 028701
- [23] Chakrabarti D, Wang Y, Wang C X, Leskovec J and Faloutsos C 2008 ACM Trans. Inf. Syst. Secur. 10 13
- [24] Restrepo J G, Ott E and Hunt B R 2006 Chaos 16 015107
- [25] Canright G S and Engø-Mosen K 2006 Comlexus 3 131
- [26] Smythe R T and Mahmoud H 1995 Theor. Probab. Math. Stat. 51 1
- [27] Erdös P and Rényi A 1960 Publ. Math. Inst. Hung. Acad. Sci. 5 17
- [28] Dorogovtsev S N, Krapivsky P L and Mendes J F F 2008 Europhys. Lett. 81 30004
- [29] Zhang Z Z, Zhou S G, Zhao S H and Guan J H 2008 J. Phys. A: Math. Theor. 41 185101
- [30] Moon J W 1974 Lond. Math. Soc. Lecture Note 13 125
- [31] Najock D and Heyde C 1982 J. Appl. Probab. 19 675

- [32] Gastwirth J 1977 Am. Stat. 31 79
- [33] Jung S, Kim S and Kahng B 2002 Phys. Rev. E 65 056101
- [34] Barabási A-L, Ravasz E and Vicsek T 2001 Physica A 299 559
- [35] Dorogovtsev S N, Goltsev A V and Mendes J F F 2002 *Phys. Rev.* E 65 066122
- [36] Comellas F, Fertin G and Raspaud A 2004 Phys. Rev. E 69 037104
- [37] Zhang Z Z, Rong L L and Zhou S G 2007 Physica A 377 329
- [38] Ravasz E and Barabási A-L 2003 Phys. Rev. E 67 026112
- [39] Nacher J C, Ueda N, Kanehisa M and Akutsu T 2005 Phys. Rev. E 71 036132
- [40] Andrade J S Jr, Herrmann H J, Andrade R F S and da Silva L R 2005 Phys. Rev. Lett. 94 018702
- [41] Zhang Z Z, Rong L L and Comellas F 2006 J. Phys. A: Math. Gen. 39 3253
- [42] Comellas F, Ozón J and Peters J G 2000 Inf. Process. Lett. 76 83
- [43] Zhang Z Z, Rong L L and Guo C H 2006 Physica A 363 567
- [44] Hinczewski M 2007 Phys. Rev. E 75 061104
- [45] Barriére L, Comellas F and Dalfó C 2006 J. Phys. A: Math. Gen. 39 11739
- [46] Zhang Z Z, Zhou S G, Fang L J, Guan J H and Zhang Y C 2007 Europhys. Lett. 79 38007
- [47] Zhang Z Z, Zhou S G, Zou T, Chen L C and Guan J H 2007 Eur. Phys. J. B 60 259
- [48] Bedogne C, Masucci A P and Rodgers G J 2008 Physica A 387 2161
- [49] Boettcher S, Gonçalves B and Guclu H 2008 J. Phys. A: Math. Theor. 41 252001
- [50] Barriére L, Comellas F, Dalfó C and Fiol M A 2008 Linear Algebr. Appl. 428 1499
- [51] Zhang Z Z, Zhou S G, Qi Y and Guan J H 2008 Eur. Phys. J. B 63 507
- [52] Goh K I, Oh E, Jeong H, Kahng B and Kim D 2002 Proc. Natl. Acad. Sci. USA 99 12583
- [53] Dorogvtsev S N, Mendes J F F and Oliveira J G 2006 Phys. Rev. E 73 056122
- [54] Watts D J and Strogatz H 1998 Nature 393 440
- [55] Newman M E J 2001 Phys. Rev. E 64 016132
- [56] Barthélemy M 2004 Eur. Phys. J. B 38 163
- [57] Szabó G, Alava M and Kertész J 2002 Phys. Rev. E 66 026101
- [58] Bollobás B and Riordan O 2004 Phys. Rev. E 69 036114
- [59] Ghima C-M, Oh E, Goh K-I, Kahng B and Kim D 2004 Eur. Phys. J. B 38 193
- [60] Zhang Z Z, Zhou S G, Chen L C, Guan J H, Fang L J and Zhang Y C 2007 Eur. Phys. J. B 59 99
- [61] Zhang Z Z, Zhou S G, Chen L C and Guan J H 2008 Eur. Phys. J. B 64 277
- [62] Barabási A-L and Albert R 1999 Science 286 509
- [63] Maslov S and Sneppen K 2002 Science 296 910
- [64] Pastor-Satorras R, Vázquez A and Vespignani A 2001 Phys. Rev. Lett. 87 258701
- [65] Vázquez A, Pastor-Satorras R and Vespignani A 2002 Phys. Rev. E 65 066130
- [66] Zhang Z Z and Zhou S G 2007 Physica A 380 621
- [67] Newman M E J 2002 Phys. Rev. Lett. 89 208701
- [68] Silvester J R 2000 Math. Gaz. 84 460